Cosmology of braneworlds *Basic principles & applications*

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TFA

Intuitevely: Our 4D spacetime is a timelike hypersurface in a 5D spacetime. The matter is confined to our spacetime \mathcal{M}_4



More precisely: *p*-brane is a *p*-dim spacelike surface in *D*-dim spacetime \mathcal{M}_D , D > p + 1

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- Closed string can move through the bulk, their spectrun contains gravitons

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- Randall–Sundrum (RS) braneworld a flat braneworld embedded into 5D AdS

Action

Analogical to 4D case

$$S = -\int \mathrm{d}^5 x \sqrt{-g^{(5)}} \left(\frac{R}{2\kappa_5^2} + \Lambda_5\right) + S_{\text{fields}} ,$$

with κ_5 and Λ_5 bulk gravitational and cosmological constants

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5D Einstein equation

$$G_{ab} \equiv R_{ab} - \frac{1}{2}g_{ab}R = -\kappa_5^2 T_{ab} + g_{ab}\Lambda_5$$

A static SO(3,1) symmetric metric

$$\mathrm{d}s^2 = \mathrm{e}^{2A(Z)}\eta_{\mu\nu}\mathrm{d}x^{\mu}\mathrm{d}x^{\nu} - \mathrm{d}Z^2$$

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$$ds^{2} = a^{2}b^{2}(dt^{2} - dZ^{2}) - a^{2}\left(\frac{dr^{2}}{1 - Kr^{2}} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta d\phi^{2}\right)$$

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Embedding of FRW braneworld

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- Consider a 3+1 timelike hypersurface \mathcal{M} in 5D bulk
- Define jump of a function f by

$$[f] = \lim_{\varepsilon \to +0} [f(Z + \epsilon) - f(Z - \epsilon)] .$$

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- Decompose stress–energy as $T_{ab} = \tau_{ab} \sigma h_{ab}$
- Projected Einstein braneworld equations are

$$G_{ab}^{(4)} = 8\pi G\tau_{ab} - \Lambda_4 h_{ab} + \kappa_5^4 \pi_{ab} - E_{ab}$$



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where

$$\pi_{ab} = \frac{1}{12}\tau\tau_{ab} - \frac{1}{4}\tau_{ac}\tau_{b}{}^{c} + \frac{1}{8}h_{ab}\tau_{cd}\tau^{cd} - \frac{1}{24}\tau^{2}h_{ab}$$

and E_{ab} is electric Weyl tensor $E_{ab} = C_{acbd} n^c n^d$



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- Main didifferences:
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 - Solution The gravitational bulk influence due to E_{ab}
 - Sianchi identity implies $\kappa_5^4 \nabla^a \pi_{ab} = \nabla^a E_{ab}$

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$$\mathrm{d}s^2 = \mathrm{e}^{-2k|Z|} \eta_{\mu\nu} \mathrm{d}x^{\mu} \mathrm{d}x^{\nu} - \mathrm{d}Z^2 \;,$$

with
$$k=\sqrt{-rac{\kappa_5^2}{6}\Lambda_5}$$





Solution Coordinate singularities at $Z = \pm \infty$









Extra dim Z compactified



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- The physical source is constitued by two branes at Z = 0 and $Z = \pi r_c$
- Intrinsic singularity on the branes
- Vanishing effective cosmological constants on the branes

Small fluctuations due to a point mass added on brane

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• Equation for $h_{\mu\nu}$ admits separation of variables, $h_{\mu\nu}(x^{\rho}, Z) = \psi(Z)\Phi(x^{\rho})$ with

$$\partial^{\mu}\partial_{\mu}\Phi(x^{\rho}) = -m^2\Phi(x^{\rho})$$

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In cylindrically symmetric case in zero-mass mode

$$\Psi(r) = -\frac{B}{r}$$

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 - Electroweak spontaneous symmetry breaking at energy scales $M_{EW} \sim 10^3 \text{ GeV}$
 - Stringy effects at scales $M_{\text{Planck}} \sim 10^{19} \text{ GeV}$
 - A huge gap between M_{Planck} and M_{EW} , 16 orders (!)

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It is sufficient $kr_c \approx 50$ in order to obtain the hierarchy

FRW braneworld

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$$H^{2} = \frac{8\pi G}{3} \left[\rho_{M} + \frac{\rho_{M}^{2}}{2\sigma} \right] + \frac{\Lambda_{4}}{3} + \frac{\mu}{a^{4}}$$
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- \square Quadratic term ρ_M

$$H^{2} = \frac{8\pi G}{6} \left[\rho_{M} + \frac{\rho_{M}^{2}}{2\sigma} \right] + \frac{\Lambda_{4}}{3} + \frac{\mu}{a^{4}}$$
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For \$\rho_M >> \sigma \sigma\$ it is \$H \sim \rho_M\$
For \$\rho_M << \sigma\$ it is \$H \sim \set \sigma_M\$ - classical scenario

According to classical cosmology, the FRW universe expands into "nothing"

According to braneworld scenario, the FRW brane expands into 5D AdS spacetime